

Possibility of Boltzmann was able to derive the entropy as Noether's charge

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Abstract

Boltzmann's virgin paper is revisited. we discuss the possibility that he derive entropy as Noether's charge and introduce Planck's constant.

1 Introduction

In recent years, Sasa and Yokokura [4] has derived entropy as a Noether's charge in a semi-static assumption. However, 150 years ago, Boltzmann is a possibility that could derive entropy as Noether charge at the time of the virgin paper 1866[1]. Then, we would like to explore the possibility that introduction of Planck's constant was possible at this point. Before the discovery of Noether's theorem[2] and Planck's constant[3] which is 30-40 years ago that these considerations performed is surprising.

2 Noether's theorem

We start to introduce Noether's theorem.

Let's consider of action ,

$$I[x] = \int_{t_2}^{t_1} L(x, \dot{x}, t) dt. \quad (2.1)$$

As a variation of this action,

$$\delta I[x] = \delta \int_{t_2}^{t_1} L(x, \dot{x}, t) dt. \quad (2.2)$$

Now, we take also variation of dt , then,

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$$\delta I[x] = \int_{t_2}^{t_1} \delta L(x, \dot{x}, t) dt - \int_{t_2}^{t_1} dL(x, \dot{x}, t) \delta t + [L(x, \dot{x}, t) \delta t]_{t_2}^{t_1}. \quad (2.3)$$

This first term become

$$\int_{t_2}^{t_1} \left(\frac{\partial L}{\partial x} \delta x + \frac{\partial L}{\partial \dot{x}} \delta \dot{x} + \frac{\partial L}{\partial t} \delta t \right) dt = \int_{t_2}^{t_1} \left(\frac{\partial L}{\partial x} - \frac{d}{dt} \frac{\partial L}{\partial \dot{x}} \right) \delta x dt + \left[\frac{\partial L}{\partial \dot{x}} \delta x \right]_{t_2}^{t_1} + \int_{t_2}^{t_1} \frac{\partial L}{\partial t} \delta t dt, \quad (2.4)$$

where

$$\delta x(t) \equiv x'(t') - x(t) = x'(t + \delta t) - x(t) = x'(t) + \dot{x}'(t) \delta t - x(t) = \delta^L x(t) + \dot{x} \delta t, \quad (2.5)$$

$$\delta^L x(t) \equiv x'(t) - x(t). \quad (2.6)$$

The second term become

$$\begin{aligned} dL(x, \dot{x}, t) \delta t &= \frac{\partial L(x, \dot{x}, t)}{\partial x} dx \delta t + \frac{\partial L(x, \dot{x}, t)}{\partial \dot{x}} d\dot{x} \delta t + \frac{\partial L(x, \dot{x}, t)}{\partial t} dt \delta t \\ &= \frac{\partial L(x, \dot{x}, t)}{\partial x} \dot{x} dt \delta t + \frac{\partial L(x, \dot{x}, t)}{\partial \dot{x}} \ddot{x} dt \delta t + \frac{\partial L(x, \dot{x}, t)}{\partial t} dt \delta t \\ &= \left(\frac{\partial L(x, \dot{x}, t)}{\partial x} - \frac{d}{dt} \frac{\partial L(x, \dot{x}, t)}{\partial \dot{x}} \right) \dot{x} dt \delta t + \frac{d}{dt} \left(\frac{\partial L(x, \dot{x}, t)}{\partial \dot{x}} \dot{x} \delta t \right) dt + \frac{\partial L(x, \dot{x}, t)}{\partial t} dt \delta t \end{aligned} \quad (2.7)$$

Summary, (2.3) is

$$\begin{aligned} \delta I[x] &= \int_{t_2}^{t_1} \left(\frac{\partial L(x, \dot{x}, t)}{\partial x} - \frac{d}{dt} \frac{\partial L(x, \dot{x}, t)}{\partial \dot{x}} \right) \delta^L x dt + \left[\frac{\partial L}{\partial \dot{x}} \delta^L x + L \delta t \right]_{t_2}^{t_1} \\ &= \int_{t_2}^{t_1} \partial_{EL} L \delta^L x dt + [p(x, \dot{x}, t) \delta x - E(x, \dot{x}, t) \delta t]_{t_2}^{t_1}, \end{aligned} \quad (2.8)$$

where

$$\partial_{EL} L \equiv \left(\frac{\partial}{\partial x} - \frac{d}{dt} \frac{\partial}{\partial \dot{x}} \right) L, \quad (2.9)$$

$$p \equiv \frac{\partial L}{\partial \dot{x}}, \quad (2.10)$$

$$E \equiv p \dot{x} - L. \quad (2.11)$$

Then, we introduce Θ by

$$\Theta \equiv p \delta x + E \delta t. \quad (2.12)$$

We get

$$\delta I[x] = \int_{t_2}^{t_1} \partial_{EL} L \delta^L x dt + [\Theta]_{t_2}^{t_1}. \quad (2.13)$$

If $\delta I[x]$ take form boundary term only,

$$\delta I[x] = [F]_{t_2}^{t_1}, \quad (2.14)$$

we get final relation,

$$- \int_{t_2}^{t_1} \partial_{EL} L \delta^L x dt = [C]_{t_2}^{t_1}, \quad (2.15)$$

and Nother's charge,

$$C \equiv \Theta - F. \quad (2.16)$$

3 Discussion on dynamical derivation of Boltzmann's entropy second law[1]

Let describe his argument with the current notation, A free particle action is

$$I[x] = \int_{t_1}^{t_2} \frac{1}{2} m \dot{x}^2 dt. \quad (3.1)$$

Variating (3.1), we get

$$\delta I[x] = \int_{t_2}^{t_1} m \dot{x} \delta \dot{x} + \int_{t_2}^{t_1} \frac{1}{2} m \dot{x}^2 \delta dt. \quad (3.2)$$

In original paper, he introduced $s = \sqrt{x^2 + y^2 + z^2}$ and calulated $d\delta s$. (3.2) is organaized like

$$\delta I[x] = - \int_{t_2}^{t_1} m \ddot{x} \delta x dt + [m \dot{x} \delta x]_{t_2}^{t_1} - \int_{t_2}^{t_1} d \left(\frac{1}{2} m \dot{x}^2 \right) \delta t + \left[\frac{1}{2} m \dot{x}^2 \delta t \right]_{t_2}^{t_1}, \quad (3.3)$$

$$\delta I[x] = \int_{t_2}^{t_1} m \ddot{x} \delta x dt - m \ddot{x} \dot{x} \delta t + \left[m \dot{x} \delta x + \frac{1}{2} m \dot{x}^2 \delta t \right]_{t_2}^{t_1}, \quad (3.4)$$

$$\delta I[x] = \int_{t_2}^{t_1} m \ddot{x} \delta^L x + \left[m \dot{x} \delta x + \frac{1}{2} m \dot{x}^2 \delta t \right]_{t_2}^{t_1}. \quad (3.5)$$

Now, we think the case of $\delta x = 0$ and the equation of motion $m \ddot{x} = 0$ is established.

Boltzmann imposed the following assumption from quasi-static adiabatic process,

$$\left[\frac{1}{2} m \dot{x}^2 \delta t \right]_{t_2}^{t_1} = \int_{t_2}^{t_1} \epsilon dt \quad (3.6)$$

By this change, $\delta I[x]$ bacome

$$\delta I[x] = [\epsilon t]_{t_2}^{t_1} = \epsilon(t_1 - t_2), \quad (3.7)$$

then, we get

$$\epsilon = \frac{\delta \int \frac{1}{2} m \dot{x}^2}{t_1 - t_2}. \quad (3.8)$$

Now, Boltzmann use the definition of temperature that was introduced at the beginning of his thesis,

$$T = \frac{\int \frac{1}{2} m \dot{x}^2}{t_1 - t_2} = \frac{I}{t_1 - t_2}. \quad (3.9)$$

This definition corresponds to Energy equality distribution law. Using (3.9), we get

$$\frac{\epsilon}{T} = \frac{\delta \int \frac{1}{2} m \dot{x}^2}{\int \frac{1}{2} m \dot{x}^2} = \delta \log I. \quad (3.10)$$

Then Boltzmann derived entropy S by

$$S = \frac{\epsilon}{T} = \beta^{-1} \epsilon = \delta \log \int_{t_2}^{t_1} \frac{1}{2} m \dot{x}^2. \quad (3.11)$$

3.1 Discussion

The change (3.6) coreponds replacement of energy to microcanonical distribution by Sasa and Yokokura [4]. Therefore, macro variable coresponding Noether's charge is

$$C = \epsilon \delta t = TS \delta t, \quad (3.12)$$

using

$$t_1 - t_2 = \delta t. \quad (3.13)$$

Then we get

$$S = \frac{\beta C}{\delta t}. \quad (3.14)$$

If we put

$$\delta t = \hbar \beta, \quad (3.15)$$

then we get entropy by Noether 's charge,

$$S = C. \tag{3.16}$$

It seems that there was a possibility that Boltzmann had anticipated the era.

References

- [1] L. Boltzmann: , Über die Mechanische Bedeutung des Zweiten Hauptsatzes der Wärme-theorie, Wiener Berichte, 53: 195–220; in WA I, paper 2.(1866)
- [2] E. Noether, Nachr. Ges. Wiss. Gottingen, 235 (1918)
- [3] Planck, Max, Ann. Phys., 309 (3): 553–63 (1901)
- [4] Sasa, Shin-ichi, and Yuki Yokokura. "Thermodynamic entropy as a Noether invariant." Physical review letters 116.14 (2016): 140601.