

Research Article

Special Relativity Temperature Transformation for a Monatomic Ideal Gas

Joseph Bevelacqua¹

1. Bevelacqua Resources, United States

A temperature transformation is derived within the scope of special relativity by examining motion in two translating inertial reference frames. The transformation is based on the four-momentum. A translating monatomic ideal gas is assumed, and the energy component of the momentum four-vector is based on an assumed Maxwell-Boltzmann distribution for the energy. The desired transformation is a complex result that depends on the temperature and momentum of the gas and the relative velocity of the inertial reference frames. In addition, the temperature transformation admits the possibility of negative temperatures.

Corresponding author: J. J. Bevelacqua, bevelresou@aol.com

1. Introduction

The temperature transformation relationship of special relativity has been reviewed extensively, but no consensus has emerged. Landsberg^[1] observed that a number of authors suggested that the temperature relationship should have the form of Eq. 1. In Eq. 1, $T(v)$ is the temperature (e.g., of a monatomic ideal gas) in a moving inertial reference frame with velocity (v), relative to the rest frame temperature $T(0)$, and b is an integer. To complete the specification of Eq. 1, γ and β are defined in Eqs. 2 and 3.

Various values for the constant b have been suggested. Without experimental data, the validity of Eq. 1 and the value of b have yet to be determined. Example values have been proposed by various authors including $b = -1$ by Planck^[2], 0 by Landsberg^[3], and $+1$ by Ott^[4]. In the Planck formulation a moving body appears to cool, while Ott's result suggests a hotter moving body. Landsberg's result suggests that temperature is an invariant quantity.

These relationships are artifacts of assumptions regarding thermodynamic properties. For example, the Planck, Landsberg, and Ott relationships result from the second law being invariant, thermodynamics being form invariant, and temperature invariance, respectively^{[5][6][7]}. However, a unique temperature relationship has yet to be determined from conventional thermodynamic considerations. Alternative formulations were suggested by Bevelacqua^{[5][6][7]}.

$$T(v) = \gamma^b T(0) \quad (1)$$

$$\gamma = \frac{1}{\sqrt{1 - \beta^2}} \quad (2)$$

$$\beta = \frac{v}{c}. \quad (3)$$

$$\begin{pmatrix} 3kT'/2c \\ p'_x \\ p'_y \\ p'_z \end{pmatrix} = \begin{pmatrix} \gamma & -\beta\gamma & 0 & 0 \\ -\beta\gamma & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} 3kT/2c \\ p_x \\ p_y \\ p_z \end{pmatrix} \quad (4)$$

$$\frac{3kT'}{2c} = \frac{3\gamma kT}{2c} - \beta\gamma p_x \quad (5)$$

$$p'_x = -\frac{3\gamma k\beta T}{2c} + \gamma p_x \quad (6)$$

$$p'_y = p_y = 0 \quad (7)$$

In^[5], the author derived a temperature relationship based on kinematic considerations that did not resemble the form summarized in (1). These results suggested that moving bodies appear to cool, but the

degree of cooling depended on the form of the equipartition theorem utilized in the formulation^[5].

A four-vector with temperature as the fourth component was utilized by Bevelacqua^[6]. The four-vector approach suggested that a moving body may exhibit any of the aforementioned temperature relationships^{[2][3][4]} depending on the nature of the four-vector and relative velocity of the inertial reference frames.

Additional work suggested that form invariance of the first and second laws of thermodynamics leads to a temperature transformation that implies that a moving body appears to cool^[7]. The constraint of the continuous transformation properties of a 4-vector with temperature as the first component also suggested that temperature does not have a simple interpretation in terms of the components of a four-vector. The difficulties associated with describing black-body radiation and temperature as a component of a four-vector has also been discussed^[7]. These issues present challenges to the concept of a continuous temperature transformation within special relativity.

Additional perspectives are provided in^{[8][9][10][11][12][13][14][15][16][17][18][19][20][21][22][23][24][25][26][27][28][29][30][31][32][33][34][35][36][37][38][39][40][41][42][43][44]}. This collection of references reinforces the fact that there is currently no consensus regarding the temperature transformation within the scope of special relativity

The four-vector approaches^{[6][7]} involved *ad hoc* assumptions. This paper utilizes a more physical approach that is based on the momentum four-vector assuming a Maxwell-Boltzmann distribution of the energy of a monatomic ideal gas. The proposed approach has well-defined assumptions that facilitate the derivation of a temperature transformation within the scope of special relativity. As such, both the assumptions and results present the possibility to be investigated experimentally.

2. Assumptions

In order to better define the proposed approach for determining a temperature transformation within the scope of special relativity, it is necessary to specify a set of assumptions. These assumptions and their validity assist in determining the viability and appropriateness of the proposed methodology. These assumptions include:

1. The proposed transformation specifies a relationship between the temperatures of a monatomic ideal gas in two inertial reference frames.

2. In the rest reference frame (K), a monatomic ideal gas, having a Maxwell-Boltzmann energy distribution, has a temperature T with associated energy $3kT/2$ where k is the Boltzmann constant^[45].
3. In the moving reference frame (K'), a monatomic ideal gas, having a Maxwell-Boltzmann energy distribution, has a temperature T' with associated energy $3kT'/2$.
4. Frame K' moves with a uniform velocity v in the positive x direction relative to a fixed frame K.
5. The K' frame initially coincides with K.
6. There is no rotational motion. The only motion is translation.

$$p'_z = p_z = 0 \quad (8)$$

$$T' = \gamma T - \frac{2c\beta\gamma}{3k} p_x \quad (9)$$

$$T' = \gamma \left(T - \frac{2v p_x}{3k} \right) \quad (10)$$

$$T' = \gamma T_{\text{eff}} \quad (11)$$

$$T_{\text{eff}} = T - \frac{2v p_x}{3k} \quad (12)$$

$$p_x > \frac{3kT}{2v} \quad (13)$$

3. Theoretical Model

The proposed temperature transformation is based on the momentum four-vector and its transformation properties. For motion of the assumed monatomic gas in the positive x direction, the momentum four-vector transforms as defined in Eq. 4. This transformation leads to the relationships summarized in Eqs. 5-8.

The desired temperature transformation for a monatomic ideal gas is obtained from Eq. 5 and is provided in Eq. 9. Eq. 9 does not have the simple transformation relationship proposed in Eq. 1. The first term in Eq. 9 resembles the Ott relationship^[4] with $b = 1$, but the second term complicates the transformation relationship. The temperature transformation result depends on the relative magnitude of these two terms.

Eq. 9 can be rewritten in the form of Eq. 1 to yield Eq. 10. Eq. 10 is rewritten in the simplified form of Eq. 11 where T_{eff} is defined in Eq. 12. However, T_{eff} is not the same temperature implied by Eq. 1.

4. Results and Discussion

Eq. 10 does not have the simple functional form of Eq. 1 with well-defined integer values for b . As such, the simple form of Eq. 1 is not obtained, but Eq. 11 provides an analogue Ott relationship.

Given the assumptions summarized in Section 2, a definitive temperature transformation does not result, but depends on a number of variables noted in Eq. 10. In addition, Eq. 10 admits the possibility of negative temperatures if the momentum p_x satisfies Eq. 13.

Although negative temperatures are a unique concept, their basis and characteristics have been addressed^{[46][47][48][49]}. The concept of negative temperature and the consequences of Eq. 10 merit additional investigation.

5. Conclusions

A temperature relationship is derived within the scope of special relativity by examining motion in two translating inertial reference frames. A monatomic ideal gas is assumed, and the energy component of the momentum four-vector is

based on an assumed Maxwell-Boltzmann energy distribution. The temperature relationship is more complex than the simple relationship of Eq. 1, and also admits the possibility of negative temperatures.

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