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Decay Characteristics of Neutron Excess Argon Nuclei

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Funding: No specific funding was received for this work.

Potential competing interests: No potential competing interests to declare.

Abstract

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In neutron star mergers, neutron excess nuclei and the r-process are important factors governing the production of heavier nuclear systems. A single-particle model evaluation of argon nuclei suggests that the heaviest $Z = 18$ nucleus will have mass 69 with filling of the $3s_{1/2}$ neutron shell. $A = 51 - 69$ argon isotopes have limited experimental half-life data, but the model predicts beta decay half-lives in the range of 0.473 – 20.3 ms. Based on previous calculations for $Z = 9 - 17$, 20, 26, and 30 systems and comparisons to the ^{51}Ar - ^{56}Ar , ^{58}Ar , and ^{60}Ar calculations, summarized in the Japanese Nuclear Data Compilation, the single-particle model results likely overestimate the half-lives of $A = 51 - 69$ neutron excess argon nuclei.

1.0 Introduction

The nucleosynthesis of heavy elements occurs by three basic processes that add protons or neutrons to a nuclear system^{1,2}. The p-process adds protons and the s- or slow process and r- or rapid process adds neutrons. Capture of protons by nuclear systems produces predominantly proton-rich nuclei that tend to decay by positron emission and electron capture^{1,2}. Neutron capture creates neutron-rich nuclei, and the resulting nuclear systems depend upon the rate of neutron addition and the beta decay rates of the residual nuclei.

In the s-process neutron capture chain, the time between successive neutron captures is sufficiently long for the product nucleus to beta decay to a stable system. Within the r-process, the time between neutron captures is too short to permit decays except for very rapid beta transitions. Therefore, the r-process must occur in an environment that has a high density of neutrons. The s-process typically occurs in red giant stars. The r-process occurs in a variety of astronomical events, including supernovae explosions and stellar mergers.

Binary neutron star or neutron star and stellar-mass black hole mergers can form a massive rotating torus around a spinning black hole¹. The matter ejected from these structures and from supernovae explosions is an important source of rapid neutron capture (r-process) nucleosynthesis¹. Fully understanding the r-process requires knowledge of the properties of neutron excess nuclei involved in creating heavy nuclear systems. Unfortunately, the majority of these

neutron excess systems have never been studied².

Closing this knowledge gap was a motivation for funding facilities for rare-isotope beams (FRIB) constructed at research facilities located around the world. These facilities are located at RIKEN (Japan)^{3,4}, GSI (Germany)⁵, and Michigan State University (US)⁶. The FRIB facilities enable a new class of experiments to determine the physical properties needed by theoretical models to determine the structure of unstable neutron excess nuclei. Theoretical studies would complement experiments that provide critical information on the unstable nuclei that must be understood in order to explain nuclear abundances observed in the universe². In particular, the study of neutron excess systems and their decay properties are significant considerations in understanding the r-process, and its importance in producing the observed elements in the universe.

The study of neutron excess systems is also important for determining nuclear decay properties, nuclear structure under extreme conditions, and nuclear reaction mechanisms. Existing theoretical models have not been extensively applied to many of these neutron excess nuclei.

This paper attempts to partially fill the void by calculating the decay properties of neutron excess systems that are important in nucleosynthesis. These theoretical studies should also assist in planning future experiments associated with neutron excess systems that are far removed from the line of stability.

Neutron excess nuclei that merit study occur throughout the periodic table²⁻⁷ including nuclei in the $Z \leq 32$ range⁷. Although neutron excess nuclei occur throughout the periodic table, this paper focuses on argon systems as part of a continuing investigation of neutron excess nuclei that are of potential astrophysical significance⁸⁻¹⁸. Previous publications addressed neutron excess calcium⁸, iron⁹, fluorine¹⁰, zinc¹¹, neon¹², sodium¹³, magnesium¹⁴, aluminum¹⁵, silicon¹⁶, phosphorous¹⁷, sulfur¹⁸, and chlorine¹⁹ systems.

The study of light nuclear systems, including argon, is important for a comprehensive astrophysical interpretation of nucleosynthesis. For example, Terasawa et al.²⁰ studied the role of light neutron-rich nuclei during r-process nucleosynthesis in supernovae. Specifically, Ref. 20 noted that light neutron excess systems can significantly affect the heavy-element abundances.

Recent studies emphasize the importance of studying argon isotopes as well as their astrophysical significance²¹⁻²⁴. These studies include both theoretical as well as experimental efforts, and provide additional data to assist in clarifying a picture of the evolution of nuclear structure with increasing neutron number

Refs. 21-24 have both theoretical nuclear physics as well as astrophysical importance in predicting the production of neutron excess argon nuclei. The continuing interest in neutron excess systems suggests the importance of evaluating argon systems considerably heavier than those investigated in Refs. 21 – 24. In particular, this paper evaluates ^{45}Ar – ^{69}Ar that span a much greater range than investigated in previous calculations.

2.0 Computational Methodology

A variety of models could be applied to the investigation of neutron excess nuclei. These vary in sophistication, but the proposed model utilizes a basic single-particle approach. This is a reasonable first step because there are uncertainties in the nuclear potential that likely are more significant than the limitations introduced by a single-particle approach.

Since the method for calculating single-particle energies in a spherically symmetric potential is well-established only salient features are provided. The model used to describe the particle plus core system represents an application of the standard method of Lukasiak and Sobiczewski²⁵ and Petrovich et. al.²⁶

The binding energy E_{NLSJ} of a particle in the field of a nuclear core is obtained by solving the radial Schrödinger Equation

$$\left[\frac{\hbar^2}{2\mu} \left(\frac{d^2}{dr^2} - \frac{L(L+1)}{r^2} \right) - E_{NLSJ} - V_{LSJ}(r) \right] U_{NLSJ}(r) = 0 \quad (1)$$

where r is the radial coordinate defining the relative motion of the nuclear core and the particle; $V_{LSJ}(r)$ is the model interaction; E_{NLSJ} is the core plus particle binding energy; $U_{NLSJ}(r)$ is the radial wave function; and L , S , and J are the orbital, spin, and total angular momentum quantum numbers, respectively. N is the radial quantum number, and μ is the reduced mass.

The method of searching for E_{NLSJ} is provided by Brown, Gunn, and Gould²⁷, and the methodology of Ref. 28 is utilized to obtain a converged solution. Refs. 8 – 19 and 26 provide a more complete description of the model, its numerical solution, and further definition of the individual terms appearing in Eq. 1.

3.0 Nuclear Interaction

Nuclear stability with respect to alpha decay, beta decay, positron decay, and electron capture is addressed using the method previously published by the author and coworkers^{8-19, 26} that is similar to the approach of Ref. 29. The single-particle level spectrum is generated using a Woods-Saxon potential. Parameters of the potential are obtained from a fit to the single-particle energy levels in ²⁰⁹Pb and ²⁰⁹Bi performed by Rost³⁰. The central potential strength of the Rost interaction³⁰ has a standard form and can be explicitly defined as

$$V_0 = 51.6 \left[1 \pm 0.73 \frac{N-Z}{A} \right] \quad (2)$$

where the upper (lower) sign applies to protons (neutrons). The remaining parameters were held constant and are given by Rost³⁰: $r_0 = 1.262$ (1.295) fm, $r_{so} = 0.908$ (1.194) fm, $a = 0.70$ (0.70) fm, and $\gamma = 17.5$ (28.2) for protons (neutrons)^{26,30}. The spin-orbit interaction strength V_{so} is related to γ by the relationship³⁰:

$$V_{so} = \frac{\gamma V_0}{180} \quad (3)$$

The scaling relationships of Eqs. 2 and 3 yield reasonable fits to observed single-particles levels in ^{120}Sn and ^{138}Ba . The pairing correction term of Blomqvist and Wahlborn³¹ is used in the calculations presented herein. The pairing correction improves the predicted energies of occupied levels in ^{120}Sn , ^{138}Ba , and ^{208}Pb ²⁶.

When applied to specific nuclei, this methodology requires modification. For example, Ray and Hodgson³² note that ^{40}Ca and ^{48}Ca require different potentials to properly fit their single-particle level structure. Schwierz, Wiedenhöver, and Volya³³ also investigated ^{40}Ca and ^{48}Ca and noted that a proper fit to the single-particle levels required a different potential for each energy level. Difficulties in the selection of an appropriate potential is an additional motivation for the utilization of a single-particle model, and was noted in studies of neutron excess calcium⁸, iron⁹, fluorine¹⁰, zinc¹¹, neon¹², sodium¹³, magnesium¹⁴, aluminum¹⁵, silicon¹⁶, phosphorous¹⁷, sulfur¹⁸, and chlorine¹⁹ nuclei. Similar issues also apply to argon systems.

In view of the results of Refs. 32 and 33, the following modification is made to obtain the chlorine potential strength (V_A):

$$V_A = 51.6\lambda \left[1 \pm 0.73 \frac{N-Z}{A} \right] [1 \pm a(A)] \text{MeV} (4)$$

where λ is a potential strength multiplier that is selected to ensure consistency with available data, and $a(A)$ is a constant that is introduced to account for the variations in potential strength with A ^{32,33}. In previous neutron excess nuclei calculations for calcium⁸, iron⁹, and zinc¹¹, a value of $\lambda = 1.0$ was utilized. A λ value of 1.5 for fluorine¹⁰, neon¹², sodium¹³, magnesium¹⁴, aluminum¹⁵, silicon¹⁶, phosphorous¹⁷, sulfur¹⁸ and chlorine¹⁹ was determined by the available experimental data³⁴⁻³⁶. Given the proximity to the $A = 18$ system, a value of $\lambda = 1.5$ is also utilized for argon. Since the paper's primary purpose is investigation of the neutron excess nuclei, determining a common $a(A)$ value for the heaviest argon systems is desirable.

The heaviest mass $A = 18$ isotope³⁴⁻³⁶ suggested experimentally is ^{50}Ar . Given the expected order of energy levels, ^{50}A would have a partially filled $2p_{3/2}$ neutron single-particle level structure. Isotopes heavier than ^{50}A would require filling of the $1f_{5/2}$, $2p_{1/2}$, $1g_{9/2}$, and $3s_{1/2}$ neutron single-particle levels. The possibility of bound argon isotopes with $A \geq 51$ is addressed in subsequent discussion.

4.0 Calculation of Half-Lives

Using Eq. 4, single-particle levels are calculated for $A \geq 45$ argon isotopes. $A \geq 45$ argon nuclei were evaluated for stability with respect to alpha decay, beta decay, positron decay, electron capture, and spontaneous fission. These calculations were performed to ensure that the nuclear structure contained no interloping states or structural defects, and that any decay modes in conflict with data were identified.

The decay modes and half-lives of $69 \geq A \geq 45$ argon isotopes are summarized in Table 1, and compared to available

data³⁴⁻³⁶ and calculations incorporated in the Japanese data compilation³⁶. The alpha decay energies are calculated using the relationship based on Ref. 37

$$Q_{\alpha} = 28.3 \text{ MeV} - 2S_n - 2S_p(5)$$

where S_n and S_p are the binding energies of the last occupied neutron and proton single-particle levels, respectively. Alpha decay half-lives can be estimated from Q_{α} using standard relationships²⁵. Fortunately, no alpha decay modes occurred in the Table 1 summary of $69 \geq A \geq 45$ argon isotope decay properties.

The beta decay half-lives are determined following the log ft methodology of Wong³⁷. Allowed (first forbidden) transition half-lives were derived using the values of $\log ft = 5$ (8). Given the uncertainties in the calculated level energies, second and higher order forbidden transitions were not determined. Positron and electron capture half-lives were determined following the approach of Ref. 25. Spontaneous fission half-lives are addressed using the methods noted in Refs. 38-51.

5.0 Model Issues

Spherical single-particle energy level calculations produce reasonable results for alpha, beta, positron, and electron capture transitions^{8-19, 29-33}, and spontaneous fission³⁸⁻⁵¹. Neutron excess argon isotopes also have the potential to decay via neutron emission modes. However, these decays have not been observed in argon³⁴⁻³⁶. The single-particle model is not the best approach for neutron emission calculations, and these decay modes are not included in this paper. Therefore, the results for the heaviest neutron excess argon nuclei only include the alpha decay, beta decay, positron decay, electron capture, and spontaneous fission modes. Except as noted previously, the single-particle model should provide reasonable results for the systems considered in the paper.

6.0 Results and Discussion

Using Eq. 4, the $a(A)$ value was varied in increments of 0.0001 to assess the applicability of the proposed model to predict the decay properties of $45 \leq A \leq 69$ argon isotopes. In view of uncertainties in the model and associated interaction, a smaller increment was not deemed to be justified.

The issues associated with fitting all calcium, iron, fluorine, zinc, neon, sodium, magnesium, aluminum, silicon, phosphorous, sulfur, and chlorine nuclei with a single potential³²⁻³³ were noted in Refs. 8-19. These considerations are also applicable to the argon systems considered in this paper.

Table 1 summarizes the complete set of $69 \geq A \geq 45$ argon isotopes considered in this paper. The $69 \geq A \geq 45$ argon isotopes fill the $1f_{7/2}$ ($^{45}\text{Ar} - ^{46}\text{Ar}$), $2p_{3/2}$ ($^{47}\text{Ar} - ^{50}\text{Ar}$), $1f_{5/2}$ ($^{51}\text{Ar} - ^{56}\text{Ar}$), $2p_{1/2}$ ($^{57}\text{Ar} - ^{58}\text{Ar}$), $1g_{9/2}$ ($^{59}\text{Ar} - ^{68}\text{Ar}$), and $3s_{1/2}$ (^{69}Ar) neutron single-particle levels. ^{50}Ar is the heaviest argon system noted in Ref. 34 - 36 that has been observed experimentally. Given the extrapolation used in formulating the single-particle potential of Eq. 4, the results become more uncertain due to the paucity of data for $A > 50$ argon isotopes. The heavier $69 \geq A \geq 45$ argon isotopes fill the $1f_{5/2}$, $2p_{1/2}$,

$1g_{9/2}$, and $3s_{1/2}$ neutron single-particle levels, and are also summarized in Table 1. These systems represent the heaviest possible neutron excess systems that could occur in the $Z=18$ system.

Table 1 Calculated Single-Particle and Experimental Decay Properties of Argon Nuclei with $45 \leq A \leq 69$

Nuclide	a(A)	Half-Life (Decay Mode)	
		Experiment ^{a,b,c}	This Work
⁴⁵ Ar	-0.0472	21.5 s (β^-) ^a	21.4 s (β^-) ^d
⁴⁶ Ar	-0.0509	8.4 s (β^-) ^a	8.38 s (β^-) ^d
⁴⁷ Ar	-0.0353	1.23 s (β^-) ^a	1.23 s (β^-) ^d
⁴⁸ Ar	-0.0281	0.48 s ^a	0.479 s (β^-) ^d
⁴⁹ Ar	-0.0127	0.170 s ^a	0.170 s (β^-) ^d
⁵⁰ Ar	-0.0010	85 ms ^b	84.9 ms (β^-) ^d
⁵¹ Ar	0.0061	e,f	20.3 ms (β^-) ^m
⁵² Ar	0.0162	e,g	14.1 ms (β^-) ^m
⁵³ Ar	0.0262	e,h	10.2 ms (β^-) ^m
⁵⁴ Ar	0.0363	e,i	7.53 ms (β^-) ^m
⁵⁵ Ar	0.0464	e,j	5.74 ms (β^-) ^m
⁵⁶ Ar	0.0565	e,k	4.43 ms (β^-) ^m
⁵⁷ Ar	0.0666	e	3.50 ms (β^-) ^m
⁵⁸ Ar	0.0766	e,l	2.81 ms (β^-) ^m
⁵⁹ Ar	0.0867	e	2.27 ms (β^-) ^m
⁶⁰ Ar	0.0968	e,n	1.87 ms (β^-) ^m
⁶¹ Ar	0.1069	e	1.56 ms (β^-) ^m
⁶² Ar	0.1170	e	1.31 ms (β^-) ^m
⁶³ Ar	0.1270	e	1.11 ms (β^-) ^m
⁶⁴ Ar	0.1371	e	0.948 ms (β^-) ^m
⁶⁵ Ar	0.1472	e	0.812 ms (β^-) ^m
⁶⁶ Ar	0.1573	e	0.703 ms (β^-) ^m
⁶⁷ Ar	0.1674	e	0.613 ms (β^-) ^m
⁶⁸ Ar	0.1774	e	0.538 ms (β^-) ^m
⁶⁹ Ar	0.1875	e	0.473 ms (β^-) ^m

^aRef. 34. ^b Ref. 35. ^c Ref. 36.
^dAllowed $1d_{3/2}(n)$ to $1d_{3/2}(p)$ beta decay transition.
^eNo data provided in Ref. 34 - 36.
^f The Japanese data compilation³⁶ notes a calculated value of 14.1 ms for ⁵¹Ar.
^g The Japanese data compilation³⁶ notes a calculated value of 8.82 ms for ⁵²Ar.
^h The Japanese data compilation³⁶ notes a calculated value of 7.09 ms for ⁵³Ar.
ⁱ The Japanese data compilation³⁶ notes a calculated value of 4.69 ms for ⁵⁴Ar.
^j The Japanese data compilation³⁶ notes a calculated value of 3.74 ms for ⁵⁵Ar.
^k The Japanese data compilation³⁶ notes a calculated value of 2.44 ms for ⁵⁶Ar.
^l The Japanese data compilation³⁶ notes a calculated value of 1.41 ms for ⁵⁸Ar.
^mAllowed $1f_{5/2}(n)$ to $1f_{7/2}(p)$ beta decay transition.
ⁿ The Japanese data compilation³⁶ notes a calculated value of 0.896 ms for ⁶⁰Ar.

The neutron excess systems summarized in Table 1 were based on an evaluation of alpha, beta, electron capture, and

positron decay modes. Spontaneous fission half-lives were also evaluated, but are larger than the aforementioned decay modes.

Other decay modes that could possibly occur in neutron excess systems (e.g., n and $2n$) are not readily evaluated using a single particle model, and were not evaluated. The results of Table 1 must be viewed with this limitation. However, since the neutron decay modes tend to be much shorter than the alpha, beta, electron capture, and positron decay modes³⁴⁻³⁶, the model results provide upper bounds on the half-lives of neutron excess argon isotopes.

6.1 $50 \geq A \geq 45$ Argon Isotopes with Experimental Half-Life Data

The $^{45}\text{Ar} - ^{46}\text{Ar}$ systems fill the $1f_{7/2}$ neutron shell. $^{45}\text{Ar} - ^{46}\text{Ar}$ are best fit with $a(A)$ values between -0.0509 and -0.0472 with an average value of about -0.0491. $^{47}\text{Ar} - ^{50}\text{Ar}$ fill the $2p_{3/2}$ neutron shell, and are best fit with $a(A)$ values between -0.0353 and -0.0010. For $^{47}\text{Ar} - ^{50}\text{Ar}$, the average $a(A)$ value is about -0.0193.

^{50}Ar is the heaviest known neutron excess argon system. There is no experimental half-life data for $A > 50$ argon systems.

The $a(A)$ values for the $^{51}\text{Ar} - ^{69}\text{Ar}$ systems are based on the decreasing lifetime trends of neutron excess silicon, phosphorous, sulfur, chlorine, and argon systems³⁴. Using the $^{45}\text{Ar} - ^{50}\text{Ar}$ values, a linear extrapolation was utilized to obtain the $a(A)$ values for the $^{51}\text{Ar} - ^{69}\text{Ar}$. The derived $a(A)$ values are listed in Table 1.

Table 1 lists the half-life of the limiting decay transition (i.e., the transition that has the shortest decay half-life). For example, ^{48}Ar has ten beta decay transitions that are possible within the scope of the aforementioned single-particle model (i.e., allowed $1d_{5/2}(n)$ to $1d_{3/2}(p)$ [2.90 s], allowed $2p_{3/2}(n)$ to $2p_{3/2}(p)$ [0.818 s], allowed $2p_{3/2}(n)$ to $2p_{1/2}(p)$ [2.84 s], first forbidden $1f_{7/2}(n)$ to $1d_{3/2}(p)$ [1.28 min], allowed $1d_{3/2}(n)$ to $1d_{3/2}(p)$ [0.479 s], allowed $1f_{7/2}(n)$ to $1f_{7/2}(p)$ [0.504 s], first forbidden $1d_{3/2}(n)$ to $1f_{7/2}(p)$ [5.36 h], first forbidden $2p_{3/2}(n)$ to $1d_{3/2}(p)$ [30.9 s], first forbidden $1d_{5/2}(n)$ to $1f_{7/2}(p)$ [78.7 yr], and allowed $1f_{7/2}(n)$ to $1f_{5/2}(p)$ [35.4 min]). For ^{48}Ar the limiting beta decay mode is the allowed $1d_{3/2}(n)$ to $1d_{3/2}(p)$ [0.479 s] transition.

As noted in Table 1, the model predicts the proper decay mode for the known $69 \geq A \geq 45$ argon³⁴⁻³⁶ systems. The results for the known nuclei, summarized in Table 1, suggest that the model predictions of the neutron excess argon systems are reasonably credible.

For nuclei filling the $1f_{7/2}$ neutron shell, model predictions for ^{45}Ar and ^{46}Ar are within about 0.5% of the experimental half-lives³⁴. ^{45}Ar and ^{46}Ar decay via beta emission through a allowed $1d_{3/2}(n)$ to $1d_{3/2}(p)$ transition. $^{47}\text{Ar} - ^{50}\text{Ar}$ decay via beta emission through allowed $1d_{3/2}(n)$ to $1d_{3/2}(p)$ transitions. Model predictions for $^{47}\text{Ar} - ^{50}\text{Ar}$ are within about 0.25% of the experimental half-lives³⁴. The calculated decay modes for $^{45}\text{Ar} - ^{50}\text{Ar}$ are in agreement with Ref. 34.

6.2 $69 \geq A \geq 51$ Argon Isotopes without Experimental Half-Life Data

The $a(A)$ values for $69 \geq A \geq 51$ argon isotopes were derived from a fit based on the half-lives of ^{45}Ar - ^{50}Ar . This approach is consistent with the $a(A)$ extrapolation methodology noted in Refs. 8 – 19. The $a(A)$ values for $69 \geq A \geq 51$ argon systems are provided in Table 1.

Table 1 also summarizes calculated single-particle decay properties of argon systems that have not yet been observed³⁴⁻³⁶. These systems are nuclei of interest in astrophysical applications¹⁻²⁴.

The existence of $69 \geq A \geq 51$ argon systems, as predicted by the proposed model, is dependent on the characteristics of the interaction of Eq. 4. Although the existence of some of these systems may be an artifact of the model interaction, their study is of critical importance in understanding the role of neutron excess argon systems in nucleosynthesis.

The ^{51}Ar - ^{56}Ar systems complete filling of the $1f_{7/2}$ neutron shell. ^{51}Ar - ^{56}Ar have beta decay half-lives that decrease from 20.3 to 4.43 ms. These systems decay through allowed $1f_{5/2}(n)$ to $1f_{7/2}(p)$ beta decay transitions. Japanese Data Compilation calculations³⁶ for ^{51}Ar - ^{56}Ar are consistent with the model results.

The ^{57}Ar and ^{58}Ar systems fill the $2p_{1/2}$ neutron shell. These systems decay through an allowed $1f_{5/2}(n)$ to $1f_{7/2}(p)$ beta decay transition. The ^{57}Ar and ^{58}Ar half-lives decrease from 3.50 to 2.81 ms, respectively. Japanese Data Compilation calculations³⁶ for ^{58}Ar are also consistent with the model results.

The ^{59}Ar - ^{68}Ar systems fill the $1g_{9/2}$ neutron shell. These systems decay through an allowed $1f_{5/2}(n)$ to $1f_{7/2}(p)$ beta decay transition. The ^{59}Ar - ^{68}Ar half-lives decrease from 2.27 to 0.538 ms.

The ^{69}Ar system partially fills the $3s_{1/2}$ neutron shell. This system decays through an allowed $1f_{5/2}(n)$ to $1f_{7/2}(p)$ beta decay transition. The ^{69}Ar half-life is 0.473 ms.

No argon systems with $A > 69$ are predicted by the model. This occurs because the $3p_{3/2}$ neutron single-particle level is the last bound neutron state, and only 51 neutrons are bound in argon systems. However, in view of the model potential uncertainties, the calculated properties of the heaviest argon systems summarized in Table 1 are not definitive.

The predicted $A = 51 - 69$ argon isotopes have no experimental data, but the model predicts beta decay half-lives in the range of 0.473 – 20.3 ms. Based on calculations in $Z = 9 - 17$, 20, 26, and 30 systems⁸⁻¹⁹ and calculations summarized in the Japanese Data Compilation³⁶, the results summarized in this paper likely overestimate the beta decay half-lives of $A = 51 - 69$ neutron excess argon nuclei. The model results are also likely to be an overestimate of the half-lives because the single-particle level calculations do not evaluate the short-lived neutron decay modes in the $A = 51 - 69$ argon nuclei.

7.0 Conclusions

Single-particle level calculations suggest that neutron excess argon isotopes terminate with ^{69}Ar and filling of the $3s_{1/2}$ neutron single-particle level. The $51 \leq A \leq 69$ argon systems have predicted beta decay half-lives in the 0.473 – 20.3 ms

range, and likely overestimate the actual half-life values.

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