

# Possible Tetraquark Explanation for the Proposed $T(2900)^{++}$ and $T(2900)^0$ Structures

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## Abstract

The recently proposed  $T_{cs\text{-}bar0}^a(2900)^{++}$  and  $T_{cs\text{-}bar0}^a(2900)^0$  structures are investigated using a first-order tetraquark mass formula. This mass relationship is based on weakly bound  $D_s^+ + \pi^+$  and  $D_s^+ + \pi^-$  meson clusters. The first-order tetraquark mass formula provides a range of mass and  $J^\pi$  predictions based on  $D_s^+$  ground and excited states.

## 1.0 Introduction

The LHCb Collaboration<sup>1,2</sup> recently reported the observation of two new tetraquarks  $T_{cs\text{-}bar0}^a(2900)^{++}$  and  $T_{cs\text{-}bar0}^a(2900)^0$ , from an analysis of  $B^+ \rightarrow D\text{-}bar^0 D_s^+ \pi^+$  and  $B^0 \rightarrow D\text{-}bar^0 D_s^+ \pi^-$  decays<sup>1,2</sup>. Refs. 1 and 2 reported these exotic states with the possible quark content of  $c\text{-}s\text{-}bar\text{-}u\text{-}d\text{-}bar$  and  $c\text{-}s\text{-}bar\text{-}d\text{-}u\text{-}bar$ . Both tetraquarks were predicted to have a mass of  $2908 \pm 11$  (statistical)  $\pm 20$  (systematic) MeV/c<sup>2</sup>. The LHCb Collaboration's data suggests a  $0^+$  assignment for these candidate states designated as  $T_{cs\text{-}bar0}^a(2900)^{++}$  and  $T_{cs\text{-}bar0}^a(2900)^0$ . The  $0^+$  value is favored over other  $J^\pi$  combinations by at least  $7.5\sigma$ .

In this paper, the first-order tetraquark mass formulas of Refs. 3 - 13 are applied to evaluate the possible mass and  $J^\pi$  values of the candidate  $T_{cs\text{-}bar0}^a(2900)^{++}$  and  $T_{cs\text{-}bar0}^a(2900)^0$  structures based on  $D_s^+ (c\text{-}s\text{-}bar) + \pi^+ (u\text{-}d\text{-}bar)$  and  $D_s^+ (c\text{-}s\text{-}bar) + \pi^- (d\text{-}u\text{-}bar)$  meson clusters.

## 2.0 Model and Formulation

Zel'dovich and Sakharov<sup>14,15</sup> proposed a semiempirical mass formula that provides a prediction of mesons and baryons in terms of effective quark masses. Within this formulation, quark wave functions are assumed to reside in their lowest 1S state. These meson mass formulas are used as the basis for deriving a first-order tetraquark mass formula. In particular, the model proposed in this paper assumes the tetraquark is partitioned into two meson clusters with the interaction between the clusters providing a minimal contribution to the tetraquark mass.

The meson mass ( $M_m$ ) formula of Refs. 3 - 13 is:

$$M_m = \delta_m + m_1 + m_2 + b_m [m_0^2 / (m_1 m_2)] \sigma_1 \cdot \sigma_2 \quad (1)$$

where  $m_1$  ( $m_2$ ) are the mass of the first (second) quark comprising the meson,  $m_0$  is the average mass of a first generation quark<sup>16,17</sup>, and the  $\sigma_i$  ( $i = 1$  and  $2$ ) are the spin vectors for the quarks incorporated into the meson. The parameters  $\delta_m$  and  $b_m$  are 40 MeV/c<sup>2</sup> and 615 MeV/c<sup>2</sup>, respectively<sup>15</sup>.

The last term in Eq. 1 represents the spin-spin interaction of the quarks and  $\sigma_1 \cdot \sigma_2$  is the scalar product of the quark

spin vectors.  $\sigma_1 \cdot \sigma_2$  has the value  $-3/4$  and  $+1/4$  for pseudoscalar and vector mesons, respectively<sup>15</sup>.

In formulating the tetraquark mass formula, effective quark masses provided by Griffiths<sup>16</sup> are utilized. These effective masses for d, u, s, c, b, and t quarks are 340, 336, 486, 1550, 4730, and 177000 MeV/c<sup>2</sup>, respectively. The effective masses are utilized in Eq. 1.

These six quarks are arranged in three generations:  $[d(-1/3), u(+2/3)]$ ,  $[s(-1/3), c(+2/3)]$ , and  $[b(-1/3), t(+2/3)]$ <sup>7</sup>. The three generations are specified by the square brackets and the quark charges [in elementary charge units (e)] are given within parentheses.

### 3.0 First-Order Mass Formula for the $T_{cs\text{-}bar0}^a(2900)^{++}$ and $T_{cs\text{-}bar0}^a(2900)^0$

The spin of a tetraquark within the first-order mass formula is determined by coupling the two meson clusters

$$\mathbf{J}^\pi = \mathbf{J}^\pi(1) \times \mathbf{L} \times \mathbf{J}^\pi(2)$$

where the first-order mass formula assumes a minimally interacting  $L=0$  configuration<sup>3-13</sup> between the meson clusters. Eq. 2 provides a primitive  $\mathbf{J}^\pi$  assignment using the possible meson clusters. The  $D_s^+$ ,  $\pi^+$ , and  $\pi^-$  have a  $0^-$  assignment. Applying Eq. 2 yields a  $0^- \times 0 \times 0^- = 0^+$  assignment in agreement with data<sup>1,2</sup>. The reader should note that there are additional excited states of the  $D_s^+$ , and these will be addressed in subsequent discussion.

The first-order mass formula used in this paper partitions the tetraquark into two meson clusters. These clusters include the  $D_s^+$  and  $\pi^+$  and  $D_s^+$  and  $\pi^-$  pseudoscalar mesons (sm). Using this structure, the tetraquark mass formula involving ground state meson clusters is assumed to have the form<sup>3-13</sup>

$$M(D_s^+ + \pi^+) = M(D_s^+ + \pi^-) = M_{sm} + M_{sm} + \Phi(3)$$

where  $\Phi$  defines the interaction between the meson clusters. Within the scope of this mass formula, the meson-meson cluster interaction is assumed to be weak and sufficiently small to be ignored. Accordingly, Eq. 3 represents the  $T_{cs\text{-}bar0}^a(2900)^{++}$  and  $T_{cs\text{-}bar0}^a(2900)^0$  structure as a quasimolecular four quark systems whose basic character is a weakly bound meson-meson system where the mesons reside in their ground states.

### 4.0 Results and Discussion

The angular momentum coupling from Eq. 2 and the first-order mass formula of Eqs. 1 and 3 are used to construct the  $T_{cs\text{-}bar0}^a(2900)^{++}$  and  $T_{cs\text{-}bar0}^a(2900)^0$  states. As noted previously, the spin and parity assignment for the  $T_{cs\text{-}bar0}^a(2900)^{++}$  and  $T_{cs\text{-}bar0}^a(2900)^0$  states is derived from Eq. 2. The resulting  $\mathbf{J}^\pi = 0^+$  assignment is in agreement with Refs. 1 and 2. As noted in Refs 3 – 13, the first-order mass formula only provides a primitive spin and parity assignment for the meson-meson cluster configuration.

Eqs. 1 and 3 lead to a predicted  $T_{cs\text{-}bar0}^a(2900)^{++}$  and  $T_{cs\text{-}bar0}^a(2900)^0$  mass of 2261 MeV/c<sup>2</sup>. This result is about 22% smaller than the experimental value<sup>1,2</sup>. Although these results are encouraging, they are based on a first-order mass formula with a number of uncertainties including the assumed quark masses<sup>16</sup>, and the magnitude of the meson-meson cluster interaction. However, the model does provide an initial description of the  $T_{cs\text{-}bar0}^a(2900)^{++}$  and  $T_{cs\text{-}bar0}^a(2900)^0$  tetraquarks that is in reasonable agreement with the experimental data<sup>1,2</sup>.

Refs. 1 and 2 note that the fit is not well described, “even adding more unobserved  $D^{**}$ ”. Given this uncertainty, the effects of excited  $D_s$  states are examined using Eq. 4

$$M(D_s^* + \pi^\pm) = M(D_s^+ + \pi^\pm) + \Delta(4)$$

where  $\Delta = M(D_s^*) - M(D_s^+)$ . The values determining  $\Delta$  are based on the Ref. 17 meson masses. The results utilizing Eq. 4 are summarized in Table 1.

Table 1			
$T_{cs\text{-}bar0}^a(2900)^{++}$ and $T_{cs\text{-}bar0}^a(2900)^0$ Tetraquark Predictions Based on Excited $D_s$ States			
Excited $D_s$ States <sup>a</sup>		First-Order Mass Formula	
State	$J^\Pi$	Mass (MeV/c <sup>2</sup> )	$J^\Pi$
$D_s^+$	$0^-$	2261	$0^+$
$D_s^{*+}$	a	2405	---
$D_{s0}^*(2317)^+$	$0^+$	2610	$0^-$
$D_{s1}^*(2460)^+$	$1^+$	2752	$1^-$
$D_{s1}^*(2536)^+$	$1^+$	2828	$1^-$
$D_{s2}^*(2573)$	$2^+$	2862	$2^-$
$D_{s1}^*(2700)^+$	$1^-$	3001	$1^+$
Experiment <sup>c</sup>		2908	$0^+$
<sup>a</sup> Ref. 17.			
<sup>b</sup> Not specified in Ref. 17.			
<sup>c</sup> Refs. 1 and 2.			

The first-order tetraquark model only predicts a  $J^\Pi = 0^+$  value in agreement with experiment if the  $D_s^+$  state is utilized in both meson clusters. Mass predictions for the  $T_{cs\text{-}bar0}^a(2900)^{++}$  and  $T_{cs\text{-}bar0}^a(2900)^0$  tetraquarks improve using the excited  $D_s$  state in one of the meson clusters, but none of the excited states yields a  $0^+$  assignment. The result closest to the experimental mass value is achieved using the  $D_{s2}^*(2573)$  in one of the meson clusters. A mass of 2862 MeV/c<sup>2</sup> is obtained using this state, which is 1.6% smaller than the measured value. However, a  $J^\Pi = 2^-$  value is predicted using the  $D_{s2}^*(2573)$  in one of the meson clusters that is not in agreement with data.

## 5.0 Conclusions

The recently proposed  $T_{cs\text{-}bar0}^a(2900)^{++}$  and  $T_{cs\text{-}bar0}^a(2900)^0$  structures<sup>1,2</sup> are investigated using a first-order tetraquark mass formula. This mass relationship is based on weakly bound the  $D_s^+$  and  $\pi^+$  and  $D_s^+$  and  $\pi^-$  meson clusters.

The assumed  $D_s^+$  and  $\pi^+$  and  $D_s^+$  and  $\pi^-$  meson cluster configurations lead to a mass result that is within about 22% of the measured  $T_{cs\text{-}bar0}^a(2900)^{++}$  and  $T_{cs\text{-}bar0}^a(2900)^0$  values<sup>1,2</sup>. The predicted first-order  $J^\Pi$  value of  $0^+$  is also in agreement with data.

Mass predictions for the  $T_{cs\text{-}bar0}^a(2900)^{++}$  and  $T_{cs\text{-}bar0}^a(2900)^0$  tetraquarks improve using the excited  $D_s$  states. The result closest to the experimental value of 2908 MeV/c<sup>2</sup> is achieved using the  $D_{s2}^*(2573)$  as one of the meson clusters. A mass of 2862 MeV/c<sup>2</sup> is obtained using this state and Eq. 4. However, a  $J^\Pi = 2^-$  value is predicted using the  $D_{s2}^*(2573)$  as one of the meson clusters that is not in agreement with data.

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